Antiprotons from dark matter ?



Pasquale Dario Serpico (Annecy, France) DarkWin, Natal, September 3rd 2019



Cosmic ray propagation & antiproton flux predictions in the "CR precision era"



Pasquale Dario Serpico (Annecy, France) DarkWin, Natal, September 3rd 2019

Broad outline

Introduction

A bit of history

A bit of history

V. Hess (outside Vienna, 1911)

A bit of history

V. Hess (outside Vienna, 1911)

positron discovery (Anderson '32)

Ist strange particle, Kaon ($\rightarrow 2\pi$, Rochester & Butler '47)

~1932-53: the first golden age of particle physics: e^+ , μ , π , strange particles K, Λ , Ξ , Σ

Including antiprotons, actually!

E. Amaldi et al., "Unusual event produced by cosmic rays," Nuovo Cimento 1 (1955), 492 (received 18/02/1955)

Unusual Event Produced by Cosmic Rays.

E. AMALDI, C. CASTAGNOLI, G. CORTINI, C. FRANZINETTI and A. MANFREDINI

Istituto di Fisica dell'Università - Roma Istituto Nazionale di Fisica Nucleare - Sezione di Roma

(ricevuto il 18 Febbraio 1955)

Summary. — The authors describe an event consisting of two stars respectively of about 5 and 1-2 GeV energy. The probable value of the number of accidental space coincidences that one expects to observe in the scanned volume, is about $4 \cdot 10^{-4}$. This value, although it d not allow us to exclude an accidental process, justifies the consideration of interpretations in terms of some physical process. Special attention is devoted to the production, capture and annihilation of a negative proton.

To be compared with Chamberlain *et al.* "Observation of Antiprotons" [@ Berkeley's Bevatron] Phys. Rev. 100, 947 (1955) **Received 24/10/1955** (Nobel Prize to Chamberlain & Segrè in 1959) although it does not allow us to exclude an accidental process, justifies [...] interpretations in terms of some physical process [...] capture and annihilation of a negative proton.

micron

top

These are secondary cosmic rays!

i.e. byproducts of astrophysical 'primary' particles hitting the upper atmosphere

Once man-made accelerator made controlled particle physics studies possible, CR field moved towards "astro" questions (1953, Bagnères de Bigorre) For an account, see James W. Cronin, 1111.5338

How is CR acceleration taking place?

In what type of objects?

Where and when?

How do CRs propagate to us?

more easily addressed by high-altitude (or above the atmosphere) detectors as the gallery in my opening slide

One thing clearly confirmed by these experiments $AMS \bar{p}/p$ results

Opportunity for IDM searches:

WIMP annihilations produce equal amounts and spectra of p and pbar's: searches much easier in pbar due to huge 'astro' background suppression!

Baseline hypothesis

pbar flux fully accounted for by rare collisions in the interstellar gas, Galactic analogous of secondaries in the atmosphere leading to particle discoveries in 1930-50

To which extent is the baseline true? Any hints for excesses attributable to "primary" contribution from DM annihilation? Rest of the talk!

The CR "Standard Model"

Sketch of the "CR standard model"

Key hypothesis

Factorized problem (differences in time and spatial scales): Sources \otimes Propagation \otimes Solar System effects (solar modulation)

Each block has a 'fiducial framework', if not a complete model

Often simplified geometry inspired by actual galactic magnetic halos

radio-contours and B-field direction of NGC 891, MW-like Galaxy

Sketch of the "standard model": sources

need to satisfy several requirements (like for artificial accelerators!)

Energetics:

need to supply 2 10⁴¹ erg/s; supernova explosions are the only class of sources providing 5-10 times more energy in the Galaxy in kinetic form (high-velocity remnants) (possibly some contributions from pulsars & stellar winds...)

Sketch of the "standard model": sources

need to satisfy several requirements (like for artificial accelerators!)

Energetics:

need to supply 2 10⁴¹ erg/s; supernova explosions are the only class of sources providing 5-10 times more energy in the Galaxy in kinetic form (high-velocity remnants) (possibly some contributions from pulsars & stellar winds...)

Mechanism for Energy Transfer:

how to transfer energy from macroscopic objects into the (microscopic) acceleration of particles?

Diffusive shock acceleration provides

I. Ind order acceleration in the large shock velocity: efficient!

II. spectral index is universal for strong (M>>I) shocks in ordinary matter

III. Spectral index value close to what inferred

$$\frac{\Delta E}{E} \simeq \frac{V_{\rm sh}}{c} \simeq 10^{-3} - 10^{-2}$$
$$E^{-\gamma} \qquad \gamma = 2 + \frac{4}{\mathcal{M}^2}$$

Sketch of the "standard model": sources

need to satisfy several requirements (like for artificial accelerators!)

Energetics:

need to supply 2 10⁴¹ erg/s; supernova explosions are the only class of sources providing 5-10 times more energy in the Galaxy in kinetic form (high-velocity remnants) (possibly some contributions from pulsars & stellar winds...)

Mechanism for Energy Transfer:

how to transfer energy from macroscopic objects into the (microscopic) acceleration of particles?

Diffusive shock acceleration provides

I. Ind order acceleration in the large shock velocity: efficient!

II. spectral index is universal for strong (M>>I) shocks in ordinary matter

III. Spectral index value close to what inferred

Confinement: need to check that the particle stays in the accelerator for the time needed to accelerate it.

Lack of (significant) E-losses: accelerating particles is useless for explaining CRs if they lose energy too quickly...

$$\frac{\Delta E}{E} \simeq \frac{V_{\rm sh}}{c} \simeq 10^{-3} - 10^{-2}$$
$$E^{-\gamma} \qquad \gamma = 2 + \frac{4}{\mathcal{M}^2}$$

Sketch of the SM: How do CRs propagate?

Charged particles deflected in B-fields, known to permeate the interstellar medium. Their "Larmor Radius" is r_L

~PeV and smaller than Galactic Sizes up to EeV.

Sketch of the SM: How do CRs propagate?

Charged particles deflected in B-fields, known to permeate the interstellar medium. Their "Larmor Radius" is r_1

$$r_L = rac{p_\perp}{Z \, e \, B} pprox rac{1 \, \mathrm{pc}}{Z} \left(rac{p_\perp}{\mathrm{PeV}/c}
ight) \left(rac{1 \mu \mathrm{G}}{B}
ight)$$

Even for protons, this distance is comparable to distance between neighboring stars up to ~PeV and smaller than Galactic Sizes up to EeV.

CRs probe thus "small-scale inhomogeneities" in the field, changing direction by what appear "random kicks", similar to Brownian motion

Macroscopically described as **diffusion (+ a drift)**

 $\frac{\partial \Phi}{\partial t} + \nabla \cdot \mathbf{J}_{\Phi} = Q$

Continuity Equation

 $\mathbf{J}_{\mathbf{\Phi}} = -D(\mathbf{x}, \Phi, \, E \ldots)
abla \Phi$ Fick's law

B⊗

Collisionless diffusion

The momentum-dependence of the diffusion depends on how large field fluctuations are at different scales (their "power spectrum")

Benchmark: "Kolmogorov" energy spectrum ~ k^{-5/3} unique dissipation rate at "inertial" scales, far from both injection (large) & viscous ones (small)

Slight complications

Add sources, energy losses... and account for the fact that the scattering centers (inhomogeneities/waves) are not static with respect to the Galactic frame.

In one dimension, CR distribution function f obeys:

Typically solved (in 1 or 2D) with numerical (GALPROP, DRAGON...) or semianalytical codes (USINE)

CR propagation for dummies

The most important effect is diffusion, especially at high-rigidity

For **stationary, homogeneous & isotropic** problems & observations at a single location, the diffusion operator can be effectively replaced by an effective "diffusive confinement" time T_{diff}

$$\frac{\partial \Phi}{\partial t} - K \nabla^2 \Phi = Q \Rightarrow \frac{\partial \Phi}{\partial t} + \frac{\Phi}{\tau_{\text{diff}}(E)} = Q$$

At steady state

$$\Phi = Q(E)\tau_{\rm diff}(E)$$

CR propagation for dummies

The most important effect is diffusion, especially at high-rigidity

For stationary, homogeneous & isotropic problems & observations at a single location, the diffusion operator can be effectively replaced by an effective "diffusive confinement" time T_{diff}

$$\frac{\partial \Phi}{\partial t} - K \nabla^2 \Phi = Q \Rightarrow \frac{\partial \Phi}{\partial t} + \frac{\partial \Phi}{\tau_0}$$

 $\frac{\Psi}{\operatorname{diff}(E)} = Q$

At steady $\Phi = Q(E)\tau_{\text{diff}}(E)$

 $\Phi_s = Q_s \tau_{\text{diff}} \propto \sigma_{p \to s} \Phi_p \tau_{\text{diff}}$ If a nucleus is not accelerated directly (primary) but only produced via primary collisions (secondary)

(more in general, convolution of σ and $\boldsymbol{\Phi}_{p}$)

- If diffusion and/or σ is E-dependent, secondaries and primaries differ in spectral shape
- Since there are no (significant) antiprotons in SNRs or the ISM, anti-p are secondaries (in the standard model); DM, as 'direct' source of anti-p, would be a primary channel!

"Fixing" the propagation parameters

Fragile nuclei such as Li, Be, B... present but in traces in stellar astrophysical environments, while in sizable fractions in CRs:

interpreted as result of spallation of "primary" nuclei, accelerated at sources (e.g. SNRs) during the CR diffusive propagation in the ISM.

Above $\sim GeV$ scale, σ has little energy dependence

"Fixing" the propagation parameters

Fragile nuclei such as Li, Be, B... present but in traces in stellar astrophysical environments, while in sizable fractions in CRs:

interpreted as result of spallation of "primary" nuclei, accelerated at sources (e.g. SNRs) during the CR diffusive propagation in the ISM.

Above $\sim GeV$ scale, σ has little energy dependence

a secondary/primary comparison yields

$$\frac{\Phi_s}{\Phi_p} \propto \tau_{\rm diff}(E) \propto 1/K(E) \propto E^{-\delta}$$

(modulo uncertainties in the x-section!)

often inferred from B/C, with typical results

$$K(R) \sim 10^{28} \div 10^{29} \left(\frac{R}{3\,{\rm GV}}\right)^{0.2 \div 0.6} {\rm cm}^2/{\rm s}$$

see e.g. Trotta, Johannesson, Moskalenko et al. ApJ 729, 106 (2011)

Sketch of the SM: Solar modulation

Result of balance between inward diffusion of Galactic CRs and outward convection and adiabatic cooling by the solar wind. Usually treated in the force-free approximation with $\varphi \sim O(100-1000)$ MV, in reality a 3D transport phenomenon in the heliosphere

Gleeson, L. J.; Axford ApJ 154, 1011 (1968)

 $J_{\oplus}(T-\varphi) = J_{\infty}(T) \frac{(T-\varphi)(T-\varphi+2m)}{T(T+2m)}$

Antiprotons

Predicting secondary antiprotons

Parameterize the injection fluxes and fit injection & propagation parameters (more or less) simultaneously to B/C, p,He...

differently from B, Li, Be... inelasticity low, hence sensitive to fluxes at significantly higher energies

Charged Cosmic Rays," JCAP 1801, 055 (2018) [1712.00002]

The AMS-02 data

M. Aguilar et al. [AMS Collaboration], "Antiproton Flux, Antiproton-to-Proton Flux Ratio, and Properties of Elementary Particle Fluxes in Primary Cosmic Rays Measured with the Alpha Magnetic Spectrometer on the International Space Station," Phys. Rev. Lett. 117, 091103 (2016)

A DM signal hidden in the data?

A.Cuoco, M. Krämer, M. Korsmeier, "Novel dark matter constraints from antiprotons in the light of AMS-02," Phys. Rev. Letters 118, 191102 (2017) [1610.03071]

"DM favored at around 4.5 sigmas"

See also

M.Y. Cui, Q.Yuan,Y.L. S.Tsai and Y.Z. Fan, "A possible dark matter annihilation signal in the AMS-02 antiproton data," Phys. Rev. Letters 118, 191101 (2017) [1610.03840]

...much smaller effect (~2.2 sigma local, I.I sigma global) found in

A. Reinert and M.W.Winkler, "A Precision Search for WIMPs with Charged Cosmic Rays," JCAP 1801, 055 (2018) [1712.00002]

2019: sequel(s)

A.Cuoco, J. Heisig, L. Klamt, M. Korsmeier, M. Krämer, "Scrutinizing the evidence for dark matter in cosmic ray antiprotons," 1903.01472

Confirmed excess, but at 3 sigma level (number of technical improvements and checks)

I. Cholis, T. Linden, D. Hooper, "A Robust Excess in the Cosmic-Ray Antiproton Spectrum: Implications for Annihilating Dark Matter" 1903.02549

Hint claimed at the 4.7 sigma level

S. J. Lin, X.J. Bi, Y. P. F.Yin, "Investigating the dark matter signal in the cosmic ray antiproton flux with the machine learning method" 1903.09545

From very significant to below 1 sigma, depending on model for cross sections and treatment of solar modulation.

Let's check! Our B/C analysis

L. Derome et al. "Fitting B/C cosmic-ray data in the AMS-02 era: a cookbook," Astron. Astrophys. 627, A158 (2019) [1904.08210]

Y. Genolini et al. "Cosmic-ray transport from AMS-02 B/C data: benchmark models and interpretation," Phys. Rev. D 99, 123028 (2019) [1904.08917]

- Takes into account a physics motivated model of correlation in AMS-02 errors.
- Nuclear cross sections parametrized... and treated as nuisance.
- Tries to avoid multiplication of parameters (probably crucial in p,He overfitting manifest in Cuoco et al.'s analyses): proves that simple power-laws work for C,O when accounting for high-rigidity spectral break in diffusion coefficient.
- Extend down to low-rigidities, modulation described with neutron monitor data and then normalization marginalized over.
- We confirm spectral break at high-rigidity, but also highlight likely presence of flattening in diffusion coefficient at low-R (wave damping?)

On error correlations

L. Derome et al. "Fitting B/C cosmic-ray data in the AMS-02 era: a cookbook," Astron. Astrophys. 627, A158 (2019) [1904.08210]

Error decomposition extracted from AMS-02 publications

• correlation lengths from educated guesses (e.g. 0 for statistical errors, 'infinite' for scale error) but for the data/Montecarlo calibration of the acceptance error, fitted with a single parameter

See 1904.08210 for a number of checks with mock data

Our B/C analysis

Y. Genolini et al. "Cosmic-ray transport from AMS-02 B/C data: benchmark models and interpretation," 1904.08917

K(R) for A/Z = 2 [kpc².Myr⁻¹]

 Good agreement with primary fluxes with simple power-law injections

 Inferred constraints e.g. on diffusion coefficient, to be reused for the secondary antiproton flux calculation

Consequences for antiprotons

M. Boudaud et al. "AMS-02 antiprotons are consistent with a secondary astrophysical origin," arXiv:1906.07119

some ingredients to pay attention to

- Contribution of different CR nuclei and targets
- Satisfactory and detailed fits the contributing nuclei (p, He, C, O...) at R>10 GV.
- Accounting for 'non-prompt' production (essentially from anti-hyperons) and uncertainty
- Accounting for the isospin violation effect & uncertainty
- Careful treatment and propagation of experimental and model uncertainties

Prediction of the antiproton flux (not a fit!)

I. IMHO, it looks remarkably close, for being an astro prediction!

2. Yet, is there an excess?

Intermediate rigidity 'broad bump'... compatible with the one found e.g. by *Cuoco et al.*?

Data-model distance usually quantified via χ^2 $x_i = \text{data}_i - \text{model}_i$ $\mathscr{C} = \text{covariance matrix}$ $\chi^2 = \sum_{ij} x_i (\mathscr{C}^{-1})_{ij} x_j$ Visual inspection of residuals via $Z_i = x_i / \sigma_i$

Wait!

- no model uncertainties, yet
- no account for correlated errors in AMS-02 data, in visual inspection!

Model Errors

Monte Carlo simulations to determine the errors (and correlations!) due to

Production XS (fits to collider data)
Transport (fit B/C)
Parent CR fluxes

Already by eye, looks less impressive, since σ_{tot} grows, lower significance

but beware of 'features by eye' (notably if correlations are present!)

Visual inspection of residuals

standard z-score

$$x_i = x_i / \sigma_i$$

What actually matters

"rotated" z-score
$$\tilde{z}_i = \tilde{x}_i / \tilde{\sigma}_i$$

where $\tilde{x}_i = U_{ij} x_j$ with U defined via $\tilde{\mathscr{C}} = U \mathscr{C} U^T$ such that $\tilde{\mathscr{C}}$ is diagonal with elements $\tilde{\mathscr{C}}_{ii} = \tilde{\sigma}_i^2$ such that $\chi^2 = \sum_i \tilde{z}_i^2$

Correlated errors in the data crucial for agreement

Correlated errors in the data crucial for agreement

Exp. errors alone bring agreement within 1 sigma, still true if most realistic treatment of model error considered

Correlated errors in the data crucial for agreement

Exp. errors alone bring agreement within 1 sigma, still true if most realistic treatment of model error considered using total errors in quadrature overestimate the uncertainties

Conclusion: AMS-02 data are consistent with a 2ary origin!

We used relatively simple propagation scenarios, and assumed no "primary" astro contribution, more realistic (hence rich) scenarios can only increase the dof's & broaden theory space".

Summary

The current CR precision era offers us new tools for DM searches (here I focused on the pbar's)

However, "Great responsibility inseparably follows from great power"

Une grande responsabilité est la suite inséparable d'un grand pouvoir

French Revolution Parliamentary Archives, Tome 64 : Du 2 au 16 mai 1793, Séance du mardi 7 mai 1793, page 287

with great power comes great responsibility

Amazing Fantasy #15 (1962)

- We argued that to analyse "features" at the level <~10 %, a number of effects must be taken into account (and some appear "irreducible errors" at the 1-2% level...)
- In particular, a bump-like antiproton excess attributable to DM seem to stem mostly from correlated errors incorrectly accounted for.

For the time being, no need for DM, but the constraints could be significant! Stay tuned...

Teaser: "naive" constraints in the literature

CR antiprotons have currently the greatest sensitivity to WIMP DM in the EW mass range

(to the extent you can trust the calculation)

A.Cuoco, M. Krämer, M. Korsmeier, "Novel dark matter constraints from antiprotons in the light of AMS-02," Phys. Rev. Letters 118, 191102 (2017) [1610.03071]

A. Reinert and M.W.Winkler, "A Precision Search for WIMPs with Charged Cosmic Rays," JCAP 1801, 055 (2018) [1712.00002]

